

## Electronic structures of the boron cage molecules $B_4H_4$ , $B_4Cl_4$ and $B_4F_4$

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*Ab initio* calculations have been performed on  $B_4H_4$ ,  $B_4Cl_4$  and  $B_4F_4$  in order to aid our understanding of the bonding in these compounds, which is presumably based on a tetrahedral boron cage. This cage has only 8 electrons and so is less than that expected on the basis of the usual framework electron counting rules. Basis sets with polarisation functions were used at the SCF, CI and CPF levels of theory to confirm that the  $T_d$  structures are indeed more stable than the  $D_{4h}$  ones. Davidson–Roby population analyses were able to show that many factors, including 3-centre 2-electron bonding and back-bonding from the ligand to the boron cage, are of importance in determining the relative stability of the three compounds, of which  $B_4Cl_4$  is the only one that has yet been observed experimentally.

**Key words:**  $B_4H_4$  —  $B_4Cl_4$  —  $B_4F_4$  — Stabilities — Geometries — Population analyses — Multi-centre bonding

### 1. Introduction

The closed deltahedral molecules  $B_4Cl_4$  (tetrachlorotetraborane(4)) [1–4], and  $(Me_3C)_4B_4$  (tetrakis (1,1-dimethylethyl) tetraborane(4)) [5], are unusual in that they are found to have only 8 electrons in the tetrahedral boron cage. According to the usual electron counting rules [6, 7], 10 valence electrons within the cage would be expected for a delocalised bonding network and 12 electrons (2 electrons for each of the 6 B–B bonds) for a localised one. Various mechanisms have been proposed to stabilise the presumably electron deficient boron cage and at the same time rationalise the non-existence of  $B_4H_4$  and  $B_4F_4$ . The most prominent amongst these involves a back-bonding process whereby valence electrons from

the ligand contribute to boron cage bonding [2, 5, 8]. Others suggest that the nature of the boron-ligand  $\sigma$  bonds or multi-centre boron-boron bonds are important factors [2, 8–10]. The difficulties involved in quantifying back-bonding effects have already been discussed with reference to tetrahedral cage bonding [10], nevertheless quantitative calculations have managed to confirm such processes; Davidson's study on BF [11] is a standard example. Similarly, the idea of multi-centre bonding is not new and it has been examined within the context of the bonding in  $B_4Cl_4$  and  $B_4F_4$ . However, it was recognised at the time that the basis sets may have been inadequate [8].

It is the aim of this paper to examine the bonding in the three compounds  $B_4H_4$ ,  $B_4Cl_4$  and  $B_4F_4$ . Our purpose was to determine the equilibrium structures within the confines of the  $T_d$  and  $D_{4h}$  point groups and analyse the differences in bonding across this series. Despite the high symmetry, previous efforts have been limited at best to SCF calculations with minimal basis sets [8–10]. We have extended these earlier calculations by performing SCF geometry optimisations with basis sets which contain polarisation functions. The influence of correlation effects has been evaluated using CI and CPF methods for  $B_4H_4$  and  $B_4F_4$  only, as size limitations restricted such calculations for  $B_4Cl_4$ .

The multi-centre contributions to the bonding in these compounds have been computed using a Davidson–Roby [11, 12] analysis, as a Mulliken [13] analysis does not give such effects and is also too sensitive to basis set variations. Such a thorough treatment should provide information as to which structures are stable, and by what means they are stabilised.

## 2. Computational methods

The basis sets employed were constructed from the Huzinaga primitive sets [14]. Two basis sets were used for each molecule, a smaller basis, denoted SB, of approximately DZ quality and a larger, denoted LB, incorporating in addition to a slightly expanded core basis, a single standard polarisation function on all atoms. In choosing our bases we hoped to accurately describe the ionic nature of the bonding and, in the case of  $B_4F_4$ , keep the basis to a size that would permit correlation calculations. Thus, we developed our LB basis for  $B_4F_4$  from a (10s, 6p) primitive set on F and a (8s, 4p) primitive set on B. The basis set details are given in Table 1.

The geometries, constrained to  $T_d$  or  $D_{4h}$  symmetry, were optimised using SCF and analytic gradient programs [15]. Especially for calculations with polarised basis sets, a new direct SCF (DSCF) program was responsible for significant improvements in computational efficiency [16]. As well as taking advantage of the molecular symmetry, it estimates the time required for the evaluation of integral batches, thus restricting storage of integrals to only those which are expensive to calculate. Convergence to the equilibrium geometry was determined by the criterion that the internal coordinate gradients were less than  $10^{-3}$  a.u. The CI and CPF calculations [17] for  $B_4H_4$  and  $B_4F_4$  were performed with the Karlsruhe version [18] of the Columbus system of program [19–21], at the SCF

**Table 1.** The basis sets for  $B_4H_4$ ,  $B_4F_4$  and  $B_4Cl_4$ 

	SB	LB
$B_4H_4$	(8s, 4p; 45s) → [4s, 2p; 2s]	(9s, 5p, 1d; 5s, 1p) → [5s, 3p, 1d; 3s, 1p]
$B_4F_4$	(8s, 4p; 8s, 4p) → [4s, 2p; 4s, 2p]	(8s, 4p, 1d; 10s, 6p, 1d) → [4s, 2p, 1d; 5s, 3p, 1d]
$B_4Cl_4$	(8s, 4p; 11s, 7p) → [4s, 2p; 6s, 4p]	(9s, 5p, 1d; 11s, 7p, 1d) → [5s, 3p, 1d; 6s, 4p, 1d]

Polarisation exponents: H( $p = 0.8$ ); B( $d = 0.5$ ); F( $d = 1.40$ ); Cl( $d = 0.65$ )

**Table 2.** Energies ( $E_h$ ) and SCF optimized geometries ( $a_0$ ) for the  $T_d$  and  $D_{4h}$  structures of  $B_4H_4$ , calculated with the SB and LB basis sets

Basis symmetry	SB		LB	
	$T_d$	$D_{4h}$	$T_d$	$D_{4h}$
$E$ SCF	-100.8655	-100.8151	-100.9479	-100.8735
CI	-101.1102	-100.9660	-101.3122	-101.1930
CPF	-101.1460	-100.9790	-101.3729	-101.2401
$r_{B-B}$	3.231	3.387	3.184	3.389
$r_{B-H}$	2.244	2.269	2.230	2.252

optimised geometry, and with a computational reduction in symmetry to the subgroups  $D_2$  and  $D_{2h}$  from  $T_d$  and  $D_{4h}$  respectively. The 1s core orbitals were kept frozen in the correlation calculations, as well as one high-lying virtual MO per BX moiety. Davidson-Roby population analyses [11, 12] were performed to partition the charge distribution into net atomic charges and shared electron numbers (SEnS) for two or multi-centre bonds [22]. All calculations were carried out on the Siemens 7881 and Cyber 205 at Karlsruhe.

### 3. Results and discussion

#### 3.1. Energetics

The calculated energies and optimised geometries for  $B_4H_4$ ,  $B_4Cl_4$  and  $B_4F_4$  are presented in Tables 2–4. With the basis set denoted by LB, which will be used

**Table 3.** SCF energies ( $E_h$ ) and optimized geometries ( $a_0$ ) for the  $T_d$  and  $D_{4h}$  structures of  $B_4Cl_4$ , calculated with the SB and LB basis sets

Basis symmetry	SB		LB	
	$T_d$	$D_{4h}$	$T_d$	$D_{4h}$
$E$ SCF	-1936.4965	-1936.4782	-1936.6585	-1936.6199
$r_{B-B}$	3.213 <sup>a</sup>	3.305	3.191 <sup>a</sup>	3.334
$r_{B-Cl}$	3.340 <sup>a</sup>	3.379	3.282 <sup>a</sup>	3.297

<sup>a</sup> Experimental values [1]:  $r_{B-B} = 3.23 \pm 0.08$  au;  $r_{B-Cl} = 3.21 \pm 0.06$  au

**Table 4.** Energies ( $E_h$ ) and SCF optimized geometries ( $a_0$ ) for the  $T_d$  and  $D_{4h}$  structures of  $B_4F_4$ , calculated with the SB and LB basis sets

Basis symmetry	SB		LB	
	$T_d$	$D_{4h}$	$T_d$	$D_{4h}$
$E$ SCF	-496.3977	-496.3813	-496.7551	-496.7134
CI			-497.7036	-497.6334
CPF			-497.9258	-497.8394
$r_{B-B}$	3.232	3.325	3.247	3.381
$r_{B-F}$	2.547	2.576	2.463	2.481

predominantly in the following discussions, all the  $T_d$  structures were found to be more stable than the corresponding  $D_{4h}$  ones by over 100 kJ mol<sup>-1</sup> (1 au  $\equiv E_h = 2625.5$  kJ mol<sup>-1</sup>) at the SCF level. The addition of polarisation functions to the basis (from SB to LB) produces additional stabilisation of the  $T_d$  structures over the  $D_{4h}$  ones.

This stabilisation of the  $T_d$  structures is further enhanced when correlation effects are included via CI or CPF. A rationalisation for this result is that a tetrahedral structure at the SCF level is more crowded electron-wise than a square-planar  $D_{4h}$  structure, and correlation corrections alleviate this crowding. At the CPF level, stabilisation of the  $T_d$  structure over the  $D_{4h}$  one is greater for  $B_4H_4$  (349 kJ mol<sup>-1</sup>) than for  $B_4F_4$  (227 kJ mol<sup>-1</sup>). Although no calculation beyond the SCF level was performed for  $B_4Cl_4$ , we would expect the tetrahedral structure for this molecule to be similarly stabilised over the square-planar structure by correlation effects.

At the SCF level the  $D_{4h}$  structures of  $B_4Cl_4$  and  $B_4F_4$  have positive electron affinities, i.e. they have an unoccupied molecular orbital with a negative orbital energy. If this is in fact the case then a change to the  $T_d$  form is most likely. It is worth noting that reduction to dianions is not favoured for the  $T_d$  structures, unlike most other compounds with the general formula  $B_nX_n$ , because the LUMO is a degenerate orbital of  $e$  symmetry and its occupation would distort the molecule.

All three  $B_4X_4$  molecules ( $X = Cl, F, H$ ),  $T_d$  and  $D_{4h}$ , were found to be stable with respect to dissociation into 4 BX monomers. A binding energy of  $\approx 1100$  kJ mol<sup>-1</sup> was computed for  $B_4H_4$  ( $T_d$ ) at the SCF and CI levels, while stability is further increased at the CPF level by  $\approx 200$  kJ mol<sup>-1</sup>. As CPF correctly accounts for size extensivity effects and CI does not, only the CPF calculations can be considered reliable when comparing BX with  $B_4X_4$ . The binding energy of  $B_4F_4$  ( $T_d$ ) was markedly lower, about 500 kJ mol<sup>-1</sup> with the SCF and CPF methods. In comparison, for the only molecule of the three that has yet been prepared,  $B_4Cl_4$  ( $T_d$ ), a binding energy of  $\approx 800$  kJ mol<sup>-1</sup> was calculated at the SCF level. These results suggest the possibility that  $B_4H_4$  and  $B_4F_4$  could also be prepared, although even in an inert matrix this would be difficult.

### 3.2. Geometries

The optimised geometries for  $B_4H_4$ ,  $B_4Cl_4$  and  $B_4F_4$  (in Tables 2–4, resp.) are the first that have been computed for these compounds with basis sets of at least double-zeta quality. Previously, only geometries obtained from experimental measurements or determined by optimisations with minimal basis sets had been used in theoretical calculations. For unusual, e.g. tetrahedral, structures this may be inadequate if subsequent computations, e.g., population analyses, are to be performed as the calculated charge distribution may not reflect the “true” situation. The calculated SCF  $T_d$  geometry of  $B_4Cl_4$  ( $r_{B-B} = 168.9$  pm,  $r_{B-Cl} = 173.7$  pm; N.B.,  $1 \text{ au} \equiv a_0 = 52.9177$  pm) is in good agreement with experiment ( $r_{B-B} = 171 \pm 4$  pm,  $r_{B-Cl} = 170 \pm 3$  pm); most of the difference could be attributed to correlation effects, which are not included in our treatment of  $B_4Cl_4$ .

In general it was found that B–X bond distances varied according to the size of X, as would be expected. In  $B_4F_4$  and  $B_4Cl_4$  the B–X distances were larger than in the corresponding monomers by  $0.080 a_0$  and  $0.028 a_0$  resp., indicating a possible reduction in the back-bonding effects which are so dominant in the BF [11] and BCl monomers. The B–B distances are about the same for all three molecules but react quite differently to the addition of polarisation functions to the basis. While there is a shortening of the BB distances in  $B_4H_4$  (by  $0.047 a_0$ ), and to a lesser extent, in  $B_4Cl_4$  (by  $0.022 a_0$ ), as the basis is enlarged from SB to LB, the corresponding length in  $B_4F_4$  is increased (by  $0.015 a_0$ ). These changes could be, in part, attributed to differences in the core bases but another explanation is possible. Polarisation functions generally cause a shrinkage of bond lengths, so some other factors are involved. A discussion of the electron distribution within the bonds, couched in terms of the shared electron numbers, can shed light on this problem.

### 3.3. Population analyses

Population analyses performed on the  $B_4X_4$  compounds and their corresponding BX monomers were of two types: a Mulliken analysis and a Davidson–Roby analysis. As a Mulliken analysis shows a marked basis set dependency, the only clear trend that could be discerned was that in the  $B_4X_4$  compounds,  $B_4F_4$  had the most ionic B–X bonds (net charges on B are 0.22, 0.07 and 0.04 for  $B_4F_4$ ,  $B_4Cl_4$  and  $B_4H_4$ , respectively).

The output from a Davidson–Roby analysis is more enlightening and gives the SENs [22], which reflect covalent bond strengths and elicit information about explicit multi-centre effects. The SENs for the  $B_4X_4$  compounds are given in Table 5 and those for BX in Table 6. The  $B_4X_4$  SENs show clearly the presence of strong bonding networks, particularly for the  $T_d$  structures and noticeably so in  $B_4H_4$  and  $B_4Cl_4$ . Strong two-centre  $\sigma$  bonds have typical values of  $SEN \approx 1.4$ , while typical values of the three-centre two-electron (3c–2e) SENs are of the order of 0.01 [24] in normal-valency compounds. The large three-centre (B–B–B) contributions (0.69, 0.51 and 0.26 for  $B_4H_4$ ,  $B_4Cl_4$  and  $B_4F_4$  resp.) give credence to the notion that bonding in the tetrahedral boron cages is stabilised by 3c–2e

**Table 5.** Shared electron numbers (SENs) for the  $T_d$  and  $D_{4h}$  structures of  $B_4H_4$ ,  $B_4Cl_4$  and  $B_4F_4$ , obtained at the SCF level with the LB basis

Molecule	Symmetry	SEN		
		B-B	B-B-B	B-X
$B_4H_4$	$D_{4h}$	2.00, 0.92 <sup>a</sup>	0.41	1.40
	$T_d$	1.87	0.69	1.46
$B_4Cl_4$	$D_{4h}$	1.63, 0.54 <sup>a</sup>	0.34	1.61
	$T_d$	1.60	0.51	1.44
$B_4F_4$	$D_{4h}$	1.61, 0.59 <sup>a</sup>	0.41	1.49
	$T_d$	1.37	0.26	1.39

<sup>a</sup> The first number refers to B-B neighbours, the second to B atoms diagonally opposed to one another

bonding, essentially extending over each face of the tetrahedron. Such a picture is similar to the bonding in  $B_2H_6$  ( $SEN(BHB) = 0.7$ ), where 3c-2e contributions dominate, and also in  $P_4$  ( $SEN(PPP) = 0.25$ ) [24].

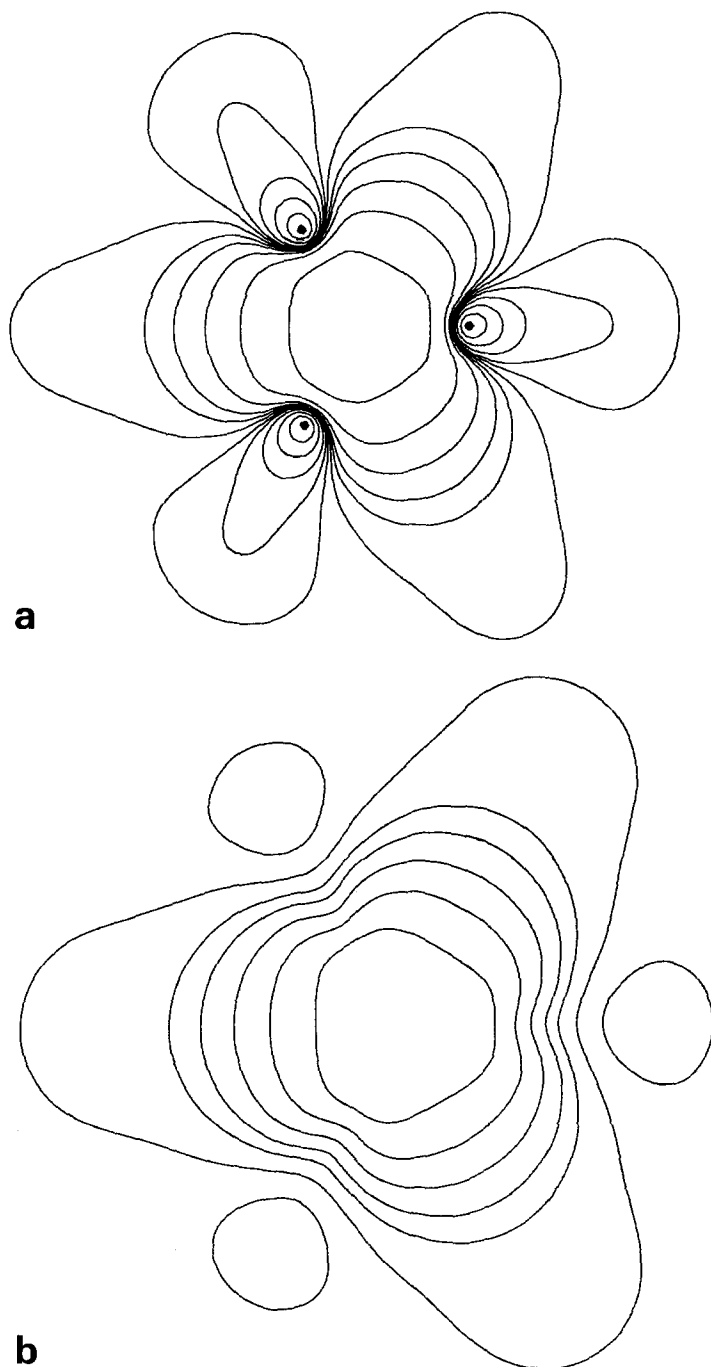
The extent of the three-centre B-B-B bonding is demonstrated more clearly if a Boys' localisation [25] is performed on the 8 bonding MOs of  $B_4H_4$ . Two sets of four equivalent Localised Molecular Orbitals (LMOs) are produced — one describing B-H bonding and the other B-B-B bonding. Figure 1a is a plot of one of the LMOs in the plane of three boron atoms and Fig. 1b is of the same LMO 0.5  $a_0$  above (external to the tetrahedron) the B-B-B plane. Both plots are typical for 3c-2e bonding. The electron density is actually found to bulge out of the plane of the three boron atoms, as is obvious if the two plots are superimposed.

The SENs computed for the BX monomers may be compared with those of the BX moieties in  $B_4X_4$ . In the monomers  $SEN(B-X)$  increases across the series BH, BCl, BF (see Table 6), while this order is reversed in the  $B_4X_4$  molecules. This observation of an apparent weakening of the B-F bond in  $B_4F_4$  compared to that of BF, together with the bond length changes discussed in the last section, indicate a decreased level of back-bonding in the B-F bonds of  $B_4F_4$ .

**Table 6.** Energies ( $E_h$ ), SCF optimized geometries ( $a_0$ ) and shared electron numbers (SENs) for BH, BCl and BF, calculated with the LB basis

	BH	BCl	BF
$E$ SCF	-25.127211	-484.089113	-124.138953
$E$ CI	-25.214071		-124.413279
$E$ CPF	-25.217561		-124.433544
$r_{B-X}$	2.310	3.254	2.383
$SEN(B-X)$	1.40	1.81 ( $p_\pi = 0.66$ )	2.00 ( $p_\pi = 0.75$ )
$r_{B-X}^a$	2.3289	3.2426	2.3860

<sup>a</sup> Experimental values [23]



**Fig. 1.** Localised molecular orbital (LMO) for  $B_4H_4$  ( $T_d$  symmetry, SCF calculation, LB basis) plotted (a) in the plane defined by 3 boron atoms and (b)  $0.05 a_0$  above (external to the tetrahedron) the plane in (a)

It is instructive to decompose the contributions to the B–B and B–X SENs into those arising from the valence MOs of  $B_4H_4$ ,  $B_4Cl_4$  and  $B_4F_4$ . This decomposition is given in Table 7. The major difference in the bonding of the three compounds is that the partly antibonding  $e$ ,  $t_1$  and  $t_2$  HOMOs of  $B_4Cl_4$  and  $B_4F_4$ , which arise from the valence  $p^\pi$  AOs (relative to the B–X bond) of Cl, F and B, are not occupied in  $B_4H_4$  because of the absence of  $p$  electrons on H. The unoccupied  $t_2$  orbital in  $B_4H_4$  (from the  $p^\pi$  AOs on B) can thus stabilise the occupied  $t_2$  orbitals (from the B–H bonds) through the coupling which is allowed through symmetry. Such a stabilisation is weaker in  $B_4Cl_4$  and  $B_4F_4$  since the  $p_\pi$  AOs on B are partly occupied, through back-bonding from the ligand, and so the occupied  $t_2$  orbitals contribute less to the two-centre SENs.

The differences between  $B_4Cl_4$  and  $B_4F_4$  are much more complex. There is competition between two mechanisms: the electron withdrawing ability of the ligand through the  $\sigma$  network, which will be greater for F than for Cl; and the ability of the ligand to supply charge to the vacant  $p$  orbitals on B (back-bonding mechanism). On comparing the BF and BCl monomers it is found that both effects are greater in BF (see Table 6,  $p_\pi = 0.75$  and  $0.66$  for BF and BCl, resp.). However, in the tetramers, back-bonding is reduced with respect to the monomers as the B  $p$  orbitals are also engaged in cage bonding. As a consequence, the net atomic charge in our Davidson–Roby analysis [22] is just slightly more positive on B in  $B_4F_4$  than in  $B_4Cl_4$  or  $B_4H_4$  ( $0.07$ ,  $-0.03$  and  $0.02$ , resp.). This reduced back-bonding is further highlighted if comparisons of the net atomic charges are made with the monomers ( $-0.21$ ,  $-0.16$  and  $0.02$  for BF, BCl and BH, resp.). Although it is difficult from the above to draw definitive conclusions it would appear that  $B_4H_4$  and  $B_4F_4$  constitute the extreme cases and  $B_4Cl_4$  lies somewhere in between, probably closer to  $B_4F_4$ .

Let us now consider the construction of  $B_4X_4$  from four BX molecules being brought in from infinity and the resulting orbital configuration ( $a_1$  and  $t_2$ ) arising

**Table 7.** Orbital energies ( $\epsilon$ ) and contributions to B–B and B–X SENs from the valence orbitals of  $B_4H_4$ ,  $B_4Cl_4$  and  $B_4F_4$  ( $T_d$  symmetry), obtained at the SCF level with the LB basis set

Orbital	$B_4H_4$			$B_4Cl_4$			$B_4F_4$		
	$\epsilon$	B–B	B–H	$\epsilon$	B–B	B–Cl	$\epsilon$	B–B	B–F
$t_2$				–0.408	0.092	–0.365	–0.442	0.096	–0.263
$t_1$				–0.498	–0.073	0.072	–0.706	–0.077	0.085
$e$				–0.514	–0.036	0.133	–0.715	–0.018	0.116
$t_2$	–0.399	0.227	0.064	–0.542	0.217	0.365	–0.724	0.121	0.271
$a_1$	–0.582	0.004	0.350	–0.618	0.031	0.155	–0.790	0.053	0.098
$t_2$	–0.618	0.338	0.993	–0.649	0.181	0.426	–0.806	–0.039	0.296
$a_1$	–0.939	1.238	0.039	–0.970	0.996	0.008	–0.966	1.145	0.017
$t_2$				–1.136	0.008	0.438	–1.646	–0.037	0.531
$a_1$				–1.145	0.136	0.189	–1.647	0.074	0.225
Total (all orbitals)		1.87	1.46		1.60	1.44		1.37	1.39



from the overlap of orbitals in the centre of the tetrahedron being formed. Implicit then in our calculations is the ability of these orbitals to rearrange so as to alleviate the increasing repulsive forces. In  $B_4H_4$ , where no back-bonding occurs, these orbitals are free to admix with the  $p_\pi$  orbitals on B ( $t_2$  symmetry) and thus move to the faces of the tetrahedron. In contrast, movement of the corresponding orbitals in  $B_4F_4$  and  $B_4Cl_4$  is hindered by the back-bonding from the ligands, even though it is less than in BF and BCl. It appears as if a compromise is required between the forces causing back-bonding effects and those forming multi-centre cage orbitals.

The results given here are generally in agreement with the previous theoretical calculations that have been performed [8–10]. Although the basis sets used here and the level of theory could be improved, the qualitative picture is unlikely to change. Davidson–Roby population analyses are known to be relatively insensitive to the nature of the basis set, however, test calculations (see also [26]) using energy-optimal, rather than standard, polarisation functions did produce small changes (reductions) in the computed SENs, particularly for the monomers, though not sufficient to affect our conclusions.

#### 4. Summary

We have recognised some important trends displayed by the  $B_4X_4$  molecules. An interplay of a number of mechanisms including the degree of multi-centre bonding, the extent of back-bonding from the ligands to the tetrahedral boron cage and the electron withdrawing capability of the ligand through the  $\sigma$  bonding network, all contribute to the stabilisation of the  $T_d$  structures over  $D_{4h}$  structures. If the compounds are ordered as follows:  $B_4H_4$ ,  $B_4Cl_4$  and  $B_4F_4$ ; then this is precisely the order of decreasing stability (with respect to dissociation into BX monomers), increasing B–B bond length, and decreasing level of bonding (as given by the shared electron numbers), both within the boron cage and to the ligands. The molecule  $B_4Cl_4$  is the only one that has been identified experimentally, but we predict that  $B_4H_4$  should also exhibit a high degree of stability, though in the liquid phase the hydrogen ligands may not provide sufficient protection from attack to the boron cage. Our prediction of reasonable stability for  $B_4H_4$  can be used to help rationalise the high stability of  $(Me_3C)_4B_4$  which does not require any back-bonding to be stable and for which the boron cage is protected by the four *t*-butyl groups. The results presented here are the most extensive yet performed on these molecules and have served to highlight the complexity of bonding in such systems; certainly the stability of molecules like  $B_4Cl_4$  cannot be attributed wholly to mechanisms such as back-bonding.

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## References

1. Atoji M, Lipscomb WN (1953) *Acta Cryst* 6:547
2. Aihara J (1983) *Bull Chem Soc Jpn* 56:335 and references therein
3. Emery SL, Morrison JA (1985) *Inorg Chem* 24:1612
4. Davan T, Morrison JA (1979) *Inorg Chem* 18:3194
5. Davan T, Morrison JA (1981) *J Chem Soc Chem Commun* 5:250
6. Wade K (1971) *J Chem Soc Chem Commun* 792; Wade K (1976) *Adv Inorg Chem Radiochem* 18:1
7. King RB, Rouvray DH (1977) *J Am Chem Soc* 99:7834
8. Hall Jr JH, Lipscomb WN (1974) *Inorg Chem* 13:710
9. Kleier DA, Bicerano J, Lipscomb WN (1980) *Inorg Chem* 19:216
10. LeBreton PR, Urano S, Shahbaz M, Emery SL, Morrison JA (1986) *J Am Chem Soc* 108:3937 and references therein
11. Davidson ER (1967) *J Chem Phys* 46:3320
12. Roby KR (1974) *Mol Phys* 27:81
13. Mulliken RS (1955) *J Chem Phys* 23:1833
14. Huzinaga S (1971) *Approximate atomic functions*. Technical Report. University of Alberta, Canada
15. Provided by Schaefer III HF and Pitzer R
16. Häser M, Ahlrichs R (1988) *J Comput Chem*, in press
17. Ahlrichs R, Scharf P, Ehrhardt C (1985) *J Chem Phys* 82:890
18. Ahlrichs R, Böhm HJ, Ehrhardt C, Scharf P, Schiffer H, Lischka H, Schindler M (1985) *J Comput Chem* 6:200
19. Dupuis M, Rys J, King HF (1976) *J Chem Phys* 65:111; Rys J, Dupuis M, King HF (1983) *J Comput Chem* 4:154
20. Pitzer RM (1973) *J Chem Phys* 58:3111
21. Lischka H, Shepard R, Brown FB, Shavitt I (1981) *Int J Quantum Chem Quantum Chem Symp* 15:91
22. Ehrhardt C, Ahlrichs R (1985) *Theor Chim Acta* 68:231
23. Huber K, Herzberg G (1979) *Molecular spectra and molecular structure. IV. Constants of diatomic molecules*. Van Nostrand Reinhold, New York
24. Ahlrichs R, Brode S, Ehrhardt C (1985) *J Am Chem Soc* 107:7260
25. Boys SF (1960) *Rev Mod Phys* 32:296
26. Heinzmann R, Ahlrichs R (1976) *Theor Chim Acta* 42:33